

## OPTICAL FIBER-BASED PLASMONIC SENSORS USING ALUMINIUM OXIDE INSULATOR

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*Abstract.* The angular and spectral interrogation methods are applied to calculation of the power loss, figure of merit and spectral and amplitude sensitivities for the TM modes in a fiber-based plasmonic sensor with four layers using a very thin aluminium oxide insulator as the interior or exterior cladding layer. The values of the effective indices calculated by using the finite element method are in agreement with an analytical method where the electromagnetic field is represented by a Bessel function of the first kind in the core region ( $\text{SiO}_2$ ), a linear combination of the Hankel functions in the gold and  $\text{Al}_2\text{O}_3$  regions, and a modified Bessel function of the second kind in the outermost region ( $\text{H}_2\text{O}$ ). For a fiber with four layers and the gold layer in the exterior part of the cladding region, the limit of detection in intensity interrogation mode ( $7.78 \times 10^{-8}$  RIU) is better than in angular interrogation mode ( $8.41 \times 10^{-6}$  RIU).

*Key words:* sensors, optical fiber, surface plasmon resonance.

### 1. INTRODUCTION

Optical fiber-based plasmonic sensors are widely used in chemical, biological, and medical monitoring [1–31]. Thus, some plasmonic biosensors [13, 19–20, 29–30] are applied for detection of human-liver tissues, of human blood groups A, B, and O and to measure the hemoglobin concentration in human blood.

The values of the effective indices calculated by using the finite element method are in agreement [10] with an analytical method where a linear combination of the Hankel functions  $H_1$  and  $H_2$  represents the field in the gold region of a fiber-based plasmonic sensor. Also, the difference between the resonant wavelengths calculated with the finite element method and with the analytical method is very small.

The interrogation (angular or spectral) method has been used for the analysis of fiber-based plasmonic sensors [22–26, 28–31].

In this paper, the angular and spectral interrogation methods are applied to an optical fiber-based plasmonic sensor when the dispersions of  $\text{SiO}_2$ , gold,  $\text{Al}_2\text{O}_3$  (insoluble in water), and distilled water ( $\text{H}_2\text{O}$ ) are considered. The power loss and amplitude sensitivity are increased at optimized thicknesses of the  $\text{Al}_2\text{O}_3$  and gold layers for the TM mode.

The values of the effective indices calculated by using the finite element method are in agreement with an analytical method where the electromagnetic field is represented by a Bessel function of the first kind in the core region ( $\text{SiO}_2$ ), a linear combination of the Hankel functions in the gold and  $\text{Al}_2\text{O}_3$  regions due to their complex values of the refractive indices, and a modified Bessel function of the second kind in the outermost region ( $\text{H}_2\text{O}$ ).

The paper is structured as follows. In Section 2 we briefly present the structure of the four-layer optical fiber-based plasmonic sensor. The angular and spectral interrogation methods are described in Section 3. We discuss in Section 4 the results of extensive numerical simulations. The summary and the discussion of the obtained results are given in Section 5.

## 2. OPTICAL FIBER-BASED PLASMONIC SENSOR

Figure 1 shows an optical fiber with four layers ( $\text{SiO}_2$ , gold,  $\text{Al}_2\text{O}_3$ , and  $\text{H}_2\text{O}$ ), where the  $\text{Al}_2\text{O}_3$  layer can be in the exterior part or in the interior part of the cladding region. The thicknesses of the  $\text{Al}_2\text{O}_3$  and gold layers are  $d_a$  and  $d_g$ , respectively.

The refractive index of the  $\text{SiO}_2$  [22],  $\text{Al}_2\text{O}_3$  [32], and distilled water [33] are calculated through a Sellmeier-type relation. The refractive index of the gold layer is calculated by the Drude model [34] as in references [7–10]. For a laser wavelength  $\lambda = 1.064 \mu\text{m}$ , the refractive indices of  $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ , gold, and  $\text{H}_2\text{O}$  are 1.449631,  $1.0818 + 9.1074i$ ,  $0.378034 + 6.21055i$  and  $1.323917 + 0.0000012193i$ , respectively.

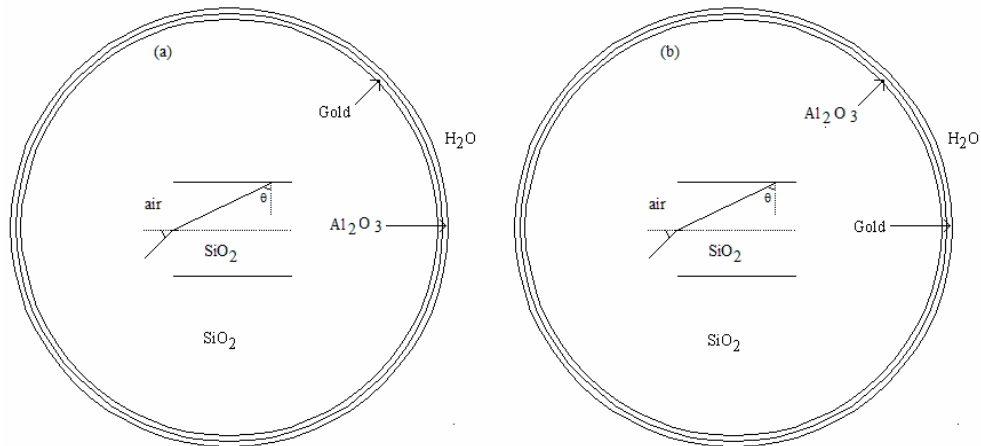


Fig. 1 – Schematic of an optical fiber with four layers ( $\text{SiO}_2$ , gold,  $\text{Al}_2\text{O}_3$  and  $\text{H}_2\text{O}$ ) in (a) and ( $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ , gold, and  $\text{H}_2\text{O}$ ) in (b). The gold layer is in interior in (a) and in exterior in (b). The light is incident from an air medium in a  $\text{SiO}_2$  core of the fiber and the angle inside the fiber is  $\theta$ .

### 3. ADAPTED ANGULAR AND SPECTRAL INTERROGATION METHODS

In the angular interrogation method [22–26, 28–31], the wavelength is kept constant and the angle of incidence is varied and a sharp dip (maximum of the power loss) appears at a resonance incidence angle  $\theta$  in the reflectivity. In the spectral interrogation method, the angle of incidence is kept constant and the wavelength is varied. A sharp dip (maximum of the power loss) appears at a wavelength  $\lambda$  in the reflectivity. The resonance incidence angle and the resonance wavelength are dependent on the refractive index of the sensing medium.

The power loss (in dB) for a *TM* mode for a Bragg fiber with  $N = 4$  layers is [25–26]:

$$PL = 10 \log \frac{1}{P}, \quad (1)$$

where  $P$  is the output power:

$$P = R^{\frac{L}{D \tan(\theta)}}, \quad (2)$$

$L$  is the sensing length,  $D$  is the fiber core diameter,  $\theta$  is the angle inside the fiber ( $L/D = 25$ ),  $R$  is the corresponding intensity reflection coefficient:

$$R = |r|^2, \quad (3)$$

and  $r$  is the amplitude reflection coefficient:

$$r = \frac{(M_{11} + M_{12}q_4)q_1 - (M_{21} + M_{22}q_4)}{(M_{11} + M_{12}q_4)q_1 + (M_{21} + M_{22}q_4)} \quad (4)$$

In the above relation we have

$$M = M_2 M_3 = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix}, \quad (5)$$

where

$$M_2 = \begin{pmatrix} \cos(\beta_2) & -\frac{i \sin(\beta_2)}{q_2} \\ -iq_2 \sin(\beta_2) & \cos(\beta_2) \end{pmatrix}, \quad M_3 = \begin{pmatrix} \cos(\beta_3) & -\frac{i \sin(\beta_3)}{q_3} \\ -iq_3 \sin(\beta_3) & \cos(\beta_3) \end{pmatrix}, \quad (6)$$

$$\beta_2 = \frac{2\pi d_2}{\lambda} \sqrt{n_2^2 - n_1^2 \sin^2(\theta)}, \quad \beta_3 = \frac{2\pi d_3}{\lambda} \sqrt{n_3^2 - n_1^2 \sin^2(\theta)}. \quad (7)$$

The real part of the effective index of the surface plasmonic wave is given by the relation:

$$\text{Re}(\beta/k) = n_1 \sin \theta = \sqrt{\frac{\varepsilon_{gr} n_{ar}^2}{\varepsilon_{gr} + n_{ar}^2}}, \quad (8)$$

where  $n_1$  is the refractive index of the core layer (SiO<sub>2</sub>),  $\theta$  is the angle inside the fiber,  $\varepsilon_{gr}$  is the real part of the gold dielectric function, and  $n_{ar}$  is the real part of the refractive index of the analyte. For a laser wavelength  $\lambda = 1.064 \mu\text{m}$ ,  $\varepsilon_{gr} = -38.428034$ ,  $n_{ar} = 1.323917$ , and  $\text{Re}(\beta/k) = 1.355184$ .

The sensing performance of the device is determined by the figure of merit (FOM):

$$\text{FOM} = \frac{\delta\theta_{res}}{\Delta n_a \delta\theta_{0.5}} \text{ or } \text{FOM} = \frac{\delta\lambda_{res}}{\Delta n_a \delta\lambda_{0.5}}, \quad (9)$$

and the maximum of the amplitude sensitivity ( $S_A$ ):

$$S_A = \frac{PL_a - PL}{\Delta n_a} \times \frac{1}{PL}, \quad \Delta n_a = 0.001 \text{ RIU}, \quad (10)$$

where  $\delta\theta_{res}$  ( $\delta\lambda_{res}$ ) is the shift in resonance angle (wavelength) corresponding to small variation in refraction index of analyte  $\Delta n_a$ ,  $\delta\theta_{0.5}$  ( $\delta\lambda_{0.5}$ ) is the resonance angular (spectral) width,  $PL_a$  is the power loss at  $n_a + \Delta n_a$ , and  $PL$  is power loss corresponding to reference analyte  $n_a$ .

For the variation of real part of the refractive index ( $\Delta n_a^r = 0.001 \text{ RIU}$ ), the angular (wavelength) resonance shift is  $\Delta\theta_{res}$  ( $\Delta\lambda_{res}$ ) and the resolution is

$$\Delta n_a^L = \frac{\Delta n_a^r}{\Delta\theta_{res}} \Delta\theta_L \text{ or } \Delta n_a^L = \frac{\Delta n_a^r}{\Delta\lambda_{res}} \Delta\lambda_L \quad (11)$$

assuming that the angular (wavelength) detection limit is  $\Delta\theta_L = 0.001^\circ$  ( $\Delta\lambda_L = 0.1 \text{ nm}$ ). For  $\Delta n_a^r = 0.001 \text{ RIU}$ , the power loss shift is  $\Delta PL_{res}$  and the resolution (limit of detection in intensity interrogation mode) is

$$\Delta n_a^L = \frac{\Delta n_a^r}{\Delta PL_{res}} \Delta P_L \quad (12)$$

assuming that the power detection limit (resolution of power measurement) is  $\Delta P_L = 0.01$  dB.

#### 4. NUMERICAL RESULTS AND DISCUSSION

Figure 2 shows the power loss *versus* the angle  $\theta$  (in radians) for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) part of the cladding region and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) part of the cladding layer at the laser wavelength  $\lambda = 1.064$   $\mu\text{m}$  for two values of the refractive index of the analyte ( $n_a$  and  $n_a + 0.001$ ). Figure 3 shows the amplitude sensitivity *versus* the angle  $\theta$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) part of the cladding region and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) part of the cladding region at the laser wavelength  $\lambda = 1.064$   $\mu\text{m}$ .

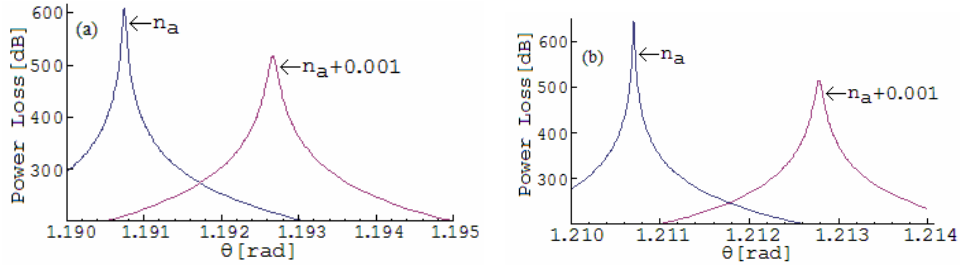


Fig. 2 – Power loss *versus* the angle  $\theta$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) in (a) and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) in (b) at the laser wavelength  $\lambda = 1.064$   $\mu\text{m}$  for two values of the refractive index of the analyte ( $n_a$  and  $n_a + 0.001$ ).

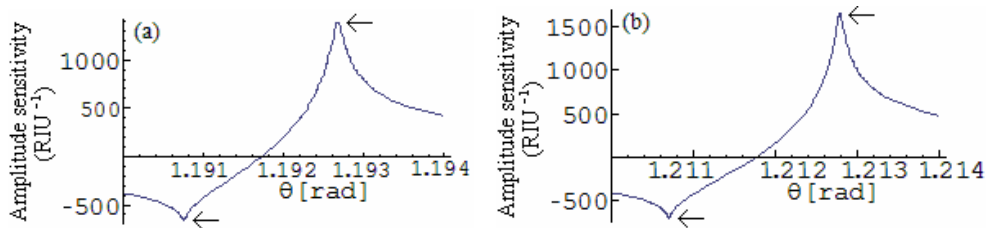


Fig. 3 – Amplitude sensitivity *versus* the angle  $\theta$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) in (a) and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) in (b) at the laser wavelength  $\lambda = 1.064$   $\mu\text{m}$ . The arrows indicate the maximum of the amplitude sensitivity.

For a TM mode, the shift towards higher angles of the maximum power loss PL for an increase  $\Delta n_a$  of the analyte refractive index by 0.001 RIU is  $\delta\theta_{\text{res}} = 0.11^\circ$  ( $0.12^\circ$ ), the full width at half maximum (FWHM) of the angular loss is  $\delta\theta_{0.5} = 0.08^\circ$  ( $0.05^\circ$ ), the figure of merit is  $\text{FOM} = 1437.2 \text{ RIU}^{-1}$  ( $2505.6 \text{ RIU}^{-1}$ ), the maximum of the amplitude sensitivity is  $S_A = 1390.6 \text{ RIU}^{-1}$  ( $1665.5 \text{ RIU}^{-1}$ ), the maximum value of the power loss is  $\text{PL} = 609.2 \text{ dB}$  ( $642.4 \text{ dB}$ ), the fixed wavelength is  $\lambda = 1.064 \mu\text{m}$  ( $1.064 \mu\text{m}$ ), and the difference between maximal amplitude sensitivity and resonant angles is  $\Delta\theta_A = 0.11^\circ$  ( $0.12^\circ$ ) when the gold layer is in the interior (exterior) part of the cladding region.

It is important that real part of the effective index  $\text{Re}(\beta/k) = n_1 \sin \theta = 1.449631 \sin(68.2249^\circ) = 1.346195$  for a fiber ( $r_1 = 1.3445 \mu\text{m}$ ,  $d_g = 19.7 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) with the gold layer in interior of the cladding region is close to  $\text{Re}(\beta/k) = 1.346353$  calculated for  $\text{HE}_{12}$  mode by using the analytical method based on the Bessel functions. The effective indices for  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes calculated with this method are  $\beta/k = 1.405147 + 0.004102i$  and  $\beta/k = 1.346353 + 0.004103i$ , respectively. For  $r_1 = 1.3445 \mu\text{m}$ , the imaginary parts of the effective indices of the  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes are very close (loss matching point 1 from Fig. 6 and Fig. 7). These values are very close to the calculated values by using the finite element method. On the other hand, the effective indices for  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes calculated with this method for a fiber ( $r_1 = 1.9296 \mu\text{m}$ ,  $d_g = 19.7 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) with the gold layer in interior of the cladding region are  $\beta/k = 1.394065 + 0.0040723i$  and  $\beta/k = 1.348518 + 0.0040721i$ , respectively. For  $r_1 = 1.9296 \mu\text{m}$ , the imaginary parts of the effective indices of the  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes are very close (loss matching point 2 from Fig. 6 and Fig. 7). Also, the real part (1.348518) of  $\beta/k$  for  $\text{TM}_{02}$  mode is close to the value (1.346195) calculated with the angular interrogation method.

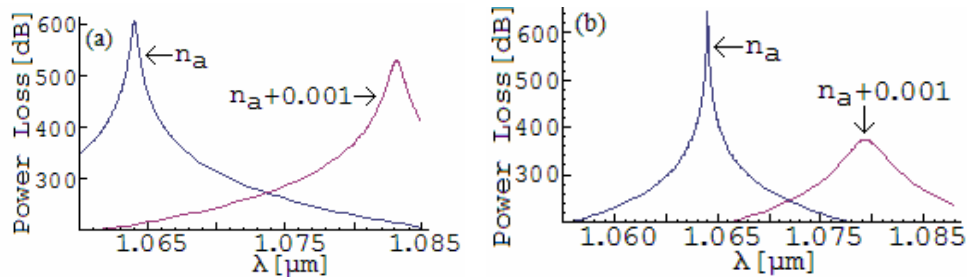


Fig. 4 – Power loss versus the wavelength  $\lambda$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) at  $\theta = 68.2249^\circ$  in (a) and for a gold layer in exterior ( $d_g = 25.9 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) at  $\theta = 69.3684^\circ$  in (b) near the wavelength  $\lambda = 1.064 \mu\text{m}$ .

Figure 4 shows the power loss versus the wavelength  $\lambda$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) at  $\theta = 68.2249^\circ$  and for a gold layer in exterior ( $d_g = 25.9 \text{ nm}$ ,  $d_a = 10 \text{ nm}$ ) at  $\theta = 69.3684^\circ$  near the wavelength  $\lambda = 1.064 \mu\text{m}$ . Figure 5 shows the amplitude sensitivity versus the wavelength  $\lambda$  for

the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) at  $\theta = 68.2249^\circ$  and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) at  $\theta = 69.3684^\circ$  near the wavelength  $\lambda = 1.064$   $\mu\text{m}$ .

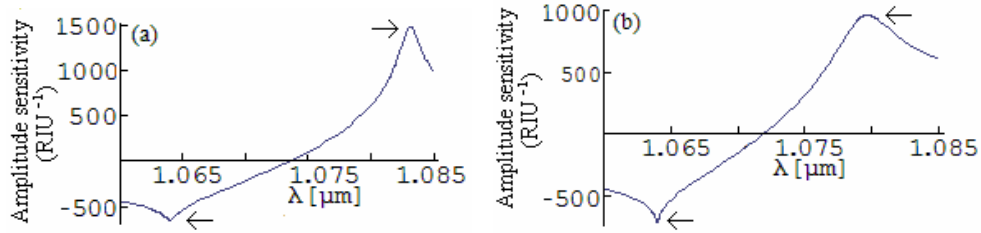


Fig. 5 – Amplitude sensitivity *versus* the wavelength  $\lambda$  for the modes TM of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) at  $\theta = 68.2249^\circ$  in (a) and for a gold layer in exterior ( $d_g = 25.9$  nm,  $d_a = 10$  nm) at  $\theta = 69.3684^\circ$  in (b) near the wavelength  $\lambda = 1.064$   $\mu\text{m}$ . The arrows indicate the maximum of the amplitude sensitivity.

For a TM mode, the shift towards higher wavelengths of the maximum power loss PL for an increase  $\Delta n_a$  of the analyte refractive index by 0.001 RIU is  $\delta \lambda_{\text{res}} = 19.11$  nm (15.34 nm), the full width at half maximum of the loss spectra is  $\delta \lambda_{0.5} = 13.26$  nm (5.91 nm), the figure of merit is  $\text{FOM} = 1440.7$   $\text{RIU}^{-1}$  (2598.1  $\text{RIU}^{-1}$ ), the maximum of the amplitude sensitivity is  $S_A = 1481.9$   $\text{RIU}^{-1}$  (962.8  $\text{RIU}^{-1}$ ), the maximum value of the power loss is  $\text{PL} = 604.6$  dB (646.1 dB), the fixed angle is  $\theta = 68.2249^\circ$  ( $69.3684^\circ$ ), and the difference between maximal amplitude sensitivity and resonant wavelengths is  $\Delta \lambda_A = 19.14$  nm (15.66 nm) when the gold layer is in interior (exterior) part of the cladding region.

Figure 6 shows the imaginary part of the effective index of the  $\text{HE}_{11}$ ,  $\text{HE}_{12}$ ,  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes as a function of the core radius  $r_1$  of the fiber with the gold in interior part of the cladding region at a fixed wavelength of 1.064  $\mu\text{m}$ . As the radius of the core layer decreases (increases), the imaginary part of the effective index of the  $\text{HE}_{11}$  ( $\text{HE}_{12}$ ) and  $\text{TM}_{01}$  ( $\text{TM}_{02}$ ) modes increases. For  $r_1 = 1.3445$   $\mu\text{m}$  there is a loss matching point between  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes and for  $r_1 = 1.9296$   $\mu\text{m}$  there is a loss matching point between  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes.

Figure 7 shows the imaginary part of the effective index of the  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes at  $r_1 = 1.3445$   $\mu\text{m}$  and for  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes at  $r_1 = 1.9296$   $\mu\text{m}$  as a function of the wavelength  $\lambda$  of the fiber with the gold in interior part of the cladding region near the wavelength of 1.064  $\mu\text{m}$ . As the wavelength decreases (increases), the imaginary part of the effective index of the  $\text{HE}_{12}$  ( $\text{HE}_{11}$ ) and  $\text{TM}_{02}$  ( $\text{TM}_{01}$ ) modes increases. For  $r_1 = 1.3445$   $\mu\text{m}$ , there is a loss matching point between  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes and for  $r_1 = 1.9296$   $\mu\text{m}$  there is a loss matching point between  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes. For  $n_a + 0.001$  the loss matching point is at  $\lambda = 1.06129$   $\mu\text{m}$  for  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes and at  $\lambda = 1.05973$   $\mu\text{m}$  for  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes.

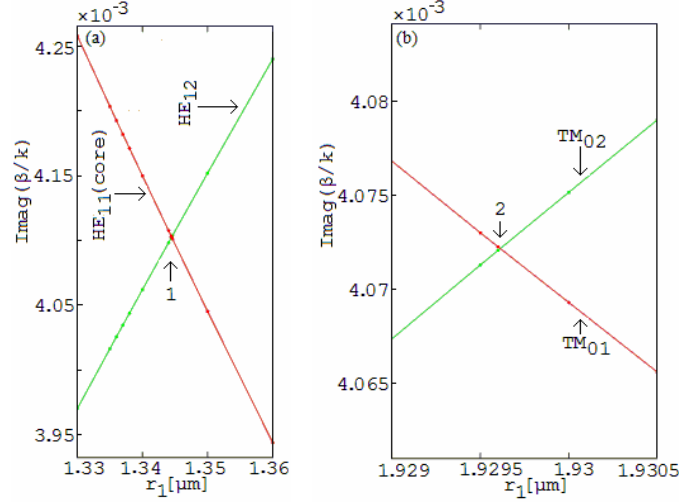


Fig. 6 – The imaginary part of the effective index  $\beta/k$  versus the core radius  $r_1$  of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) for  $HE_{11}$  core (red) and  $HE_{12}$  (green) modes (a) near  $r_1 = 1.3445$   $\mu\text{m}$  (loss matching point 1) and for  $TM_{01}$  (red) and  $TM_{02}$  (green) modes (b) near  $r_1 = 1.9296$   $\mu\text{m}$  (loss matching point 2) (Color online).

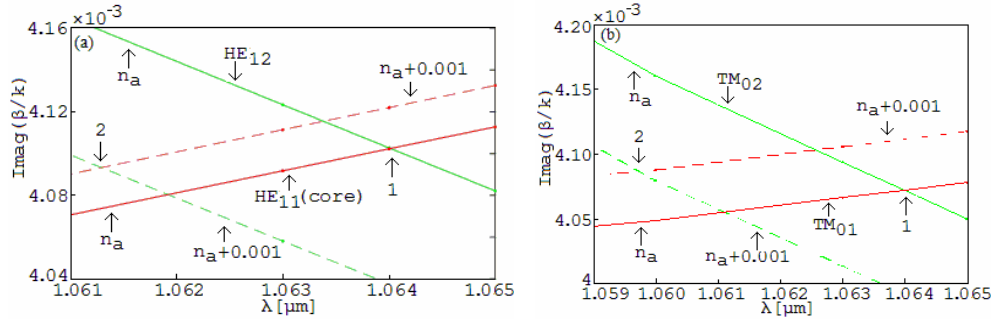


Fig. 7 – The imaginary part of the effective index  $\beta/k$  versus the wavelength  $\lambda$  of a fiber with the gold layer in interior ( $d_g = 19.7$  nm,  $d_a = 10$  nm) for  $HE_{11}$  core (red) and  $HE_{12}$  (green) modes (a) with  $r_1 = 1.3445$   $\mu\text{m}$  and for  $TM_{01}$  (red) and  $TM_{02}$  (green) modes (b) with  $r_1 = 1.9296$   $\mu\text{m}$  for two values of the refractive index of the analyte ( $n_a$  with the loss matching point 1 and  $n_a + 0.001$  with the loss matching point 2) (Color online).

Table 1 shows the shift  $\delta\theta_{\text{res}}$  ( $\delta\lambda_{\text{res}}$ ) towards higher angles (wavelengths) of the maximum power loss PL for an increase  $\Delta n_a$  of the analyte refractive index by 0.001 RIU, the full width at half maximum (FWHM) of the angular loss  $\delta\theta_{0.5}$  (loss spectra  $\delta\lambda_{0.5}$ ), the figure of merit FOM, the maximum of the amplitude sensitivity  $S_A$ , the maximum value of the power loss PL, the wavelength  $\lambda$ , the angle inside the fiber  $\theta$ , and the difference  $\Delta\theta_A$  ( $\Delta\lambda_A$ ) between maximal amplitude sensitivity and resonant angles (wavelengths) for a TM mode of a fiber when the gold layer is in the interior part of the cladding region.



Table 1

Values of  $\delta\theta_{\text{res}}$  [°],  $\delta\lambda_{\text{res}}$  [nm],  $\delta\theta_{0.5}$  [°],  $\delta\lambda_{0.5}$  [nm], FOM [RIU<sup>-1</sup>],  $S_A$  [RIU<sup>-1</sup>], PL [dB],  $\lambda$  [μm],  $\theta$  [°],  $\Delta\theta_A$  [°] and  $\Delta\lambda_A$  [nm] for a TM mode of an optical fiber with four layers with the thicknesses  $d_g$  [nm] and  $d_a$  [nm] of the gold and Al<sub>2</sub>O<sub>3</sub> layers, respectively. The gold layer is in the interior of the cladding region

$d_g; d_a$	$\delta\theta_{\text{res}}$ $\delta\lambda_{\text{res}}$	$\delta\theta_{0.5}$ $\delta\lambda_{0.5}$	FOM	$S_A$	PL	$\lambda$ $\theta$	$\Delta\theta_A$ $\Delta\lambda_A$
19.7; 10	0.11 19.11	0.08 13.26	1437.2 1440.7	1390.6 1481.9	609.2 604.6	1.064 68.2249	0.11 19.14

Table 2

Values of  $\delta\theta_{\text{res}}$  [°],  $\delta\lambda_{\text{res}}$  [nm],  $\delta\theta_{0.5}$  [°],  $\delta\lambda_{0.5}$  [nm], FOM [RIU<sup>-1</sup>],  $S_A$  [RIU<sup>-1</sup>], PL [dB],  $\lambda$  [μm],  $\theta$  [°],  $\Delta\theta_A$  [°], and  $\Delta\lambda_A$  [nm] for a TM mode of an optical fiber with four layers with the thicknesses  $d_g$  [nm] and  $d_a$  [nm] of the gold and Al<sub>2</sub>O<sub>3</sub> layers, respectively. The gold layer is in the exterior of the cladding layer. The last line is for an optical fiber with three layers without Al<sub>2</sub>O<sub>3</sub> layer

$d_a; d_g$	$\delta\theta_{\text{res}}$ $\delta\lambda_{\text{res}}$	$\delta\theta_{0.5}$ $\delta\lambda_{0.5}$	FOM	$S_A$	PL	$\lambda$ $\theta$	$\Delta\theta_A$ $\Delta\lambda_A$
10; 25.9	0.12 15.34	0.05 5.91	2505.6 2598.1	1665.5 962.8	642.4 646.1	1.064 69.3684	0.12 15.66
0; 44.4	0.12 16.42	0.06 8.02	2011.9 2045.8	1747.4 1124.0	577.8 578.6	1.064 69.4685	0.12 16.7

Table 2 shows the same characteristics as in Table 1 for a TM mode of a fiber when the gold layer is in the exterior part of the cladding region. The last line of this Table is for an optical fiber with three layers without Al<sub>2</sub>O<sub>3</sub> layer when the optimized thickness of the gold layer is increased at 44.4 nm.

For a TM mode of a plasmonic fiber without an Al<sub>2</sub>O<sub>3</sub> layer, the optimized thickness of the gold layer is  $d_g = 44.4$  nm and the sensing characteristics are given in the last line of the Table 2. It is important that the real part of the effective index  $\text{Re}(\beta/k) = n_1 \sin\theta = 1.449631 \sin(69.4685^\circ) = 1.357550$  for the fiber ( $r_1 = 1.3468$  μm,  $d_g = 44.4$  nm,  $d_a = 0$  nm) is close to  $\text{Re}(\beta/k) = 1.359846$  calculated for HE<sub>12</sub> mode by using the analytical method based on the Bessel functions. The effective indices for HE<sub>11</sub> and HE<sub>12</sub> modes calculated with this method are  $\beta/k = 1.405063 + 0.00305421i$  and  $\beta/k = 1.359846 + 0.00305374i$ , respectively. For  $r_1 = 1.3468$  μm, the imaginary parts of the effective indices of the HE<sub>11</sub> and HE<sub>12</sub> modes are very close (loss matching point). These values are very close to the calculated values by using the finite element method. On the other hand, the effective indices for TM<sub>01</sub> and TM<sub>02</sub> modes calculated with this method for the fiber ( $r_1 = 1.9610$  μm,  $d_g = 44.4$  nm,  $d_a = 0$  nm) are  $\beta/k = 1.396111 + 0.00297159i$  and  $\beta/k = 1.361327 + 0.00297136i$ , respectively. For  $r_1 = 1.9610$  μm, the imaginary parts of the effective indices of the TM<sub>01</sub> and TM<sub>02</sub> modes are very close (loss matching point). Also, the real part (1.361327) of  $\beta/k$  for TM<sub>02</sub> mode is close to the value (1.357550) calculated with the angular interrogation method.

For an angular detection limit of  $0.001^\circ$ , the sensor angular resolution ( $8.41 \times 10^{-6}$  RIU) is smaller for a fiber with four layers when the gold layer is in the exterior part of the cladding region. For the same fiber structure and a power detection limit of 0.01 dB, the sensor power resolution in intensity interrogation mode is  $7.78 \times 10^{-8}$  RIU. Thus, the limit of detection in the intensity interrogation mode is better and more useful than in the angular interrogation mode.

## 5. CONCLUSIONS

The angular and spectral interrogation methods show that for an optical fiber-based plasmonic sensor with a thin  $\text{Al}_2\text{O}_3$  insulator layer, the best figure of merit and the maximum of the amplitude sensitivity in the infrared region near the laser wavelength  $\lambda = 1.064 \mu\text{m}$  are obtained when the gold layer is in the exterior part of the cladding region of the fiber for the distilled water as the analyte layer.

For some values of the core ( $\text{SiO}_2$ ) radius there is a loss matching point between  $\text{HE}_{11}$  and  $\text{HE}_{12}$  modes and between  $\text{TM}_{01}$  and  $\text{TM}_{02}$  modes when the imaginary parts of the effective indices of the  $\text{HE}_{11}$  ( $\text{TM}_{01}$ ) and  $\text{HE}_{12}$  ( $\text{TM}_{02}$ ) modes calculated with an analytical method based on the Bessel functions are very close.

For a fiber with four layers and the gold layer in the exterior part of the cladding region, the limit of detection ( $7.78 \times 10^{-8}$  RIU) in the intensity interrogation mode is better than in the angular interrogation mode ( $8.41 \times 10^{-6}$  RIU).

It is important to note that the difference  $\Delta\theta_A$  ( $\Delta\lambda_A$ ) between maximal amplitude sensitivity and resonant angles (wavelengths) are very close to the shift  $\delta\theta_{\text{res}}$  ( $\delta\lambda_{\text{res}}$ ) towards higher angles (wavelengths) of the maximum power loss PL for an increase  $\Delta n_a$  of the analyte refractive index by 0.001 RIU.

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